

# Mass attenuation coefficient

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We all know the so-called “linear” attenuation coefficient. Here I set the symbol  $\lambda$  for that. Let  $x$  to be the symbol for the thickness of homogeneous layers, where the beam transmits perpendicular with the intensity  $I_0$ . After the layer  $x$  we have the intensity

$$I = I_0 \exp(-\lambda x). \quad (1)$$

This exponentially decreasing function of attenuation has first found by Pierre Bouguer (1698–1758) [1]. In his figure [1, Fig. 2] we can numerate the layers  $i = 1, 2, \dots$ . He found  $\log I_{i-1} - \log I_i = c(x_i - x_{i-1})$ , which agrees with (1),  $c$  is a constant. The logarithmus of intensity is a linear function of the thickness of the transmitted layer.

The  $\lambda$  is not the basic quantity of attenuation. The basic is the mass attenuation. Let us set the symbol  $\mu$  for the mass attenuation coefficient. This  $\mu$  is important in practice.

Let us suppose the beam of the intensity  $I$  has the cross area  $A$ . It goes through a layer  $dx$  with density  $\rho$  (here of one-element matter), which has the atomic mass  $M$  and cross section  $\sigma$ . There the change of intensity is

$$dI = -I \frac{\frac{\rho}{M} N dx A \sigma}{A} = -\mu \rho I dx, \quad (2)$$

$N$  is Avogadro's number.  $\frac{\rho}{M}$  is the number of moles in volume unite.  $\frac{\rho}{M} N$  is the number of atoms in

volume unite. Then  $\frac{\rho}{M} N A dx$  is the number of atoms in the differential volume  $dV = A dx$ . The

cross section  $\sigma$  is the interaction (reaction or scattering) area for an atom. The numerator in the middle of (2) is the sum of the areas where the interaction occurs. Its relation to  $A$  is the probability of the interaction.

$$\mu = \frac{N \sigma}{M} \quad (3)$$

is the mass attenuation coefficient. It is an important radiation quantity.

For the molecule, which has  $n$  elements, an element  $k$  times, we have

$$\mu = N \sum_{i=1}^n \frac{k_i \sigma_i}{M_i} \quad (4)$$

$\mu$  can also be written for a mixture of elements and molecules. For both  $\sigma$  and  $\mu$  of photons NIST has [tabulation](#) [2]. There the mass attenuation coefficient  $\mu$  is in the unit  $\text{cm}^2/\text{g}$ .

In neutron physics the mass attenuation coefficient is called the macroscopic cross section [3]. The cross section  $\sigma$  then you find in the [tabulation](#) [4].

The integration of (2) over a homogenous layer  $x$  gives the intensity

$$I = I_0 \exp(-\mu \rho x), \quad (5)$$

when  $I_0$  is the intensity before the layer.

Gardner and Ely [5] use the symbol  $\mu$  as in (5) for the attenuation.

In the XIV International Symposium on Radiation Physics, 2018 in Cordoba in Argentina, the discrepancy about symbols of attenuation coefficients still was alive. In the encyclopaedia [6] the different  $\mu$ 's are discussed.

If we have  $n$  homogeneous layers then

$$I = I_0 \exp\left(-\sum_{i=1}^n \mu_i \rho_i x_i\right), \quad (6)$$

and if the density varies in  $x$ -direction, then

$$I = I_0 \exp\left(-\mu \int_0^x \rho(x) dx\right). \quad (7)$$

But when  $x$  is large then build-up considerations are needed, [7] for gamma photons.

We have  $\lambda = \mu \rho$  as the relation between linear attenuation coefficient and mass attenuation coefficient in homogeneous matter with the density  $\rho$ .

When one determines  $\lambda$ , then also  $\mu$  and its accuracy should be determined.

I think (1) should be called the absorption law of Bouguer.

## References

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- [4] Special Issue on Nuclear Reaction Data (eds. Pavel Obložinský, Boris Pritychenko), [Nuclear Data Sheets 148](#), 1-420, February 2018.
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